NANO LETTERS

Janus Solid—Liquid Interface Enabling Ultrahigh Charging and Discharging Rate for Advanced Lithium-Ion Batteries

Jiaxin Zheng,[†] Yuyang Hou,^{#,‡} Yandong Duan,[†] Xiaohe Song,[†] Yi Wei,[†] Tongchao Liu,[†] Jiangtao Hu,[†] Hua Guo,[†] Zengqing Zhuo,[†] Lili Liu,[‡] Zheng Chang,[‡] Xiaowei Wang,[‡] Danylo Zherebetskyy,[§] Yanyan Fang,^{||} Yuan Lin,^{||} Kang Xu,[⊥] Lin-Wang Wang,[§] Yuping Wu,^{*,#,‡} and Feng Pan^{*,†}

[†]School of Advanced Materials, Peking University, Shenzhen Graduate School, Shenzhen 518055, People's Republic of China

[#]College of Energy, Nanjing Tech University, Nanjing 211816, Jiangsu Province, People's Republic of China

[‡]New Energy and Materials Laboratory (NEML), Department of Chemistry & Shanghai Key Laboratory of Molecular Catalysis and Innovative Materials, Fudan University, Shanghai 200433, People's Republic of China

[§]Materials Science Division, Lawrence Berkeley National Laboratory, Berkely, California 94720, United States

^{||}Beijing National Laboratory for Molecular Sciences, Key Laboratory of Photochemistry, Institute of Chemistry, Chinese Academy of Sciences, Beijing 100190, People's Republic of China

¹Electrochemistry Branch, U.S. Army Research Laboratory, Adelphi, Maryland 20783, United States

Supporting Information

ABSTRACT: LiFePO₄ has long been held as one of the most promising battery cathode for its high energy storage capacity. Meanwhile, although extensive studies have been conducted on the interfacial chemistries in Li-ion batteries,¹⁻³ little is known on the atomic level about the solid–liquid interface of LiFePO₄/electrolyte. Here, we report battery cathode consisted with nanosized LiFePO₄ particles in aqueous electrolyte with an high charging and discharging rate of 600 C (3600/600 = 6 s charge time, 1 C = 170 mAh g⁻¹) reaching



C (3600/600 = 6 s charge time, 1 C = 170 mAh g⁻¹) reaching 72 mAh g⁻¹ energy storage (42% of the theoretical capacity). By contrast, the accessible capacity sharply decreases to 20 mAh g⁻¹ at 200 C in organic electrolyte. After a comprehensive electrochemistry tests and ab initio calculations of the LiFePO₄-H₂O and LiFePO₄-EC (ethylene carbonate) systems, we identified the transient formation of a Janus hydrated interface in the LiFePO₄-H₂O and LiFePO₄-EC (ethylene carbonate) systems, we identified the transient formation of a Janus hydrated interface in the LiFePO₄-H₂O molecules, forming a *half-solid* (LiFePO₄) and *half-liquid* (H₂O) amphiphilic coordination environment that eases the Li desolvation process near the surface, which makes a fast Li-ion transport across the solid/liquid interfaces possible.

KEYWORDS: LiFePO₄, rate performance, aqueous electrolyte, organic electrolyte, solid-liquid interface, ab initio calculations

harging/discharging rate (CDR) is a key battery parameter that dictates how fast energy can be harnessed or released and is critical for applications such as vehicleelectrification and renewable energy grids.^{4,5} Although supercapacitors and redox-flow batteries exhibit certain potentials for these applications, 5-7 their capacities are often limited or costs are too high.⁸ On the other hand, the high energy and power densities and excellent cycling stability of lithium-ion batteries (LIBs) have made them the most promising candidate. As a safer LIB cathode chemistry, LiFePO₄ has been intensively studied in the past decade,⁹ whose CDR was found to be surprisingly decoupled from its low intrinsic ionic and electronic conductivities. Previous work has identified that, with nanosized particles, the limiting step is no longer the Li⁺intercalation rate within the particles but rather the transports of Li^+ and electrons to or from them.^{10–12} As a commonly used and highly effective practice, carbon coating has successfully addressed the electronic contact among electrode active ingredients,¹³ whereas recent experiment and theory suggested that Li⁺-transport across the cathode/electrolyte interface might play the dictating role for the LiFePO₄ rate performance.^{14–17} Gaberscek and co-workers suggested that the "ionic coating" around LiFePO₄ might matter more than electronic coating due to the much clumsier ionic movements,¹⁴ and they believe that, at sufficiently small particle sizes, the rate-limiting step should be the interfacial reaction rather than transport within the bulk.

On the other hand, electrolytes also exercise a critical control on the rate performance,^{18,19} determining not only Li⁺-diffusion within the bulk electrolyte, but also the transport process at solid/electrolyte interface. In recent years aqueous lithium battery has been intensively explored as an alternative to its nonaqueous counterparts, mainly because of the intrinsic safe

Received:June 16, 2015Revised:August 20, 2015Published:August 25, 2015





Figure 1. Morphology of the prepared 45 nm LiFePO₄ (LFP) nanoparticles and its electrochemical properties. (a) Scan electron micrograph (SEM) of the morphology. (b) Transmission electron micrograph (TEM): large percentage of the surface is (010) faced. (c) and (d) Charge and discharge curves at different current density ($1 C = 170 \text{ mA g}^{-1}$) between -0.2 and 0.75 V (vs SCE) in 0.5 M Li₂SO₄ aqueous electrolyte (c) and in an organic electrolyte (1 M LiClO₄ electrolyte in a mixture of ethylene carbonate and dimethyl carbonate) (d).

nature brought by the water-based electrolytes.^{20–25} However, many fundamental questions regarding how LIB chemistries behave in aqueous and nonaqueous media remains unanswered¹⁶ (Table S2 in Supporting Information), which includes: how differently do organic and aqueous electrolyte interact with electrode surfaces on atomic level, which factor governs the kinetics of Li⁺-intercalation/deintercalation processes at the interfaces dissecting solid electrodes from electrolytes, and what would a "perfect" interface be in order for Li⁺-transport to occur at ultrahigh rate with minimum resistance?

Here, we report our efforts in addressing these questions by comparing CDRs of LiFePO₄ in both aqueous and nonaqueous electrolytes. With cathode consisting of nanosized LiFePO₄ in 0.5 M Li₂SO₄ aqueous solution, a rate of 600 C (1 C = 170 mAh g⁻¹) could be achieved for both charging and discharging while still retaining a capacity of 72.0 mAh g⁻¹ (42% of the theoretical capacity). By contrast, the accessible capacity sharply decreases to 20 mAh g⁻¹ at 200 C in organic electrolyte. We also found that this high rate capability comes with minimum interfacial resistance and excellent cycling stability in aqueous electrolyte, that is, after 400 cycles at 10 C, the loss of capacity is only 5% (less than 15% after 1000 cycles). Combination of these merits makes aqueous LIB using such nanosized LiFePO₄ a promising candidate for applications requiring high capacity and high CDR, and with additional

benefits such as nonflammability, low cost, and small environmental footprint. All imparted from the aqueous electrolytes, render particular suitability for large-scale energy storage systems (>10⁶ Wh), where concerns over the prohibitive high cost of production, operation and maintenance, potential safety hazard, and environmental disaster far outweigh any advantages in energy densities.

Results and Discussion. Different crystallite sizes of $LiFePO_4$ (45 and 101 nm) were prepared by reflux route in ethylene glycol solution under atmospheric pressure. LiFePO₄ samples were mixed with 20 wt % of glucose and then carbonized at 650 °C for 6 h in Ar atmosphere. The obtained LiFePO₄/C composites were single LiFePO₄ phase as determined by X-ray diffraction (XRD) (Figure S1c in Supporting Information). The SEM images of LiFePO₄/C are shown in Figure 1a and Figure S1a in Supporting Information. The samples consist of monodispersed nanocrystalline particles with platelet shape, and the high-resolution transmission electron microscopy (HRTEM) images (Figure 1b and Figure S1b in Supporting Information) reveal well-crystallized nanoparticles of 45 and 101 nm. The patterns of fast Fourier transform of HRTEM displayed in the inset of Figure 1b and Figure S1b in Supporting Information indicate that the 45 nm LiFePO₄ crystal has a major exposure of (010) facets, whereas the 101 nm LiFePO₄ grows transversely along the (010)direction.



Figure 2. Kinetic properties of the electrochemical redox reactions. (a) Cyclic voltammograms (CVs) for the prepared 45 nm LiFePO₄ (LFP) nanoparticles in 0.5 M Li₂SO₄ aqueous electrolyte using Pt wire as the counter electrode and Ag/AgCl as the reference electrode, and 1 mol L⁻¹ LiClO₄ organic electrolyte (a mixture of ethylene carbonate and dimethyl carbonate) at the scan rate of 20 mV s⁻¹. Note that separations between cathodic and anodic peaks in aqueous (~0.1 V) are much smaller than that in nonaqueous (~0.5 V) electrolyte. Electrochemical kinetics numerical simulations are adopted to fit the redox peaks in the CVs (green lines, section S1 in Supporting Information), where electrochemical kinetics process model was adopted while taking into account the anodic and cathodic peak shifts upon the change of scanning rate. (b) Change of mass during the CV scan measured by EQCM in charge process and discharge process (inset), respectively. M_A and M_O denote the change of mass induced by every molar electrolyte and the organic electrolyte, respectively.

The voltage profiles as functions of capacity (the degree of intercalations) at varying CDRs are described in Figure 1c and d in 0.5 M Li₂SO₄ aqueous electrolyte and in 1 M LiClO₄ in ethylene carbonate (EC) and dimethyl carbonate (DMC), respectively. Note that ultrahigh charging and discharging curves in the aqueous and organic electrolytes are symmetric (Figure 1c and d), and there are voltage platforms (close to 3.5 V vs Li⁺/Li) in the charge and discharge curves (Figure 1c and d), which is a typical behavior for LiFePO₄ and can be clearly seen up to 600 C in aqueous electrolyte and up to 120 C (1 C = 170 mAh g^{-1}) in organic electrolyte. The drop of the voltage of the platform with the increasing charge and discharge rate is due to the increasing polarization. We can see that at ultrahigh charging rates of 600 C, a capacity of 72.0 mAh g^{-1} (42% of the theoretical capacity) was retained in the aqueous electrolyte (Figure 1c and Supporting Information Figure S2). Compared with the LiFePO₄ nanoparticles prepared similarly but evaluated in organic electrolytes, where the accessible capacity sharply decreases to 20 mAh g^{-1} at 200 C (Figure 1d), the rate capability in aqueous electrolyte is apparently much superior. Here, we also measured the capacity of the carbon (acetylene black) under the same experimental conditions and at different current densities. As shown in Figure S3 in Supporting Information, the acetylene black exhibits very low capacity both in aqueous and organic electrolytes at high current rates, so the capacity of carbon could be ignored in the working electrode (mixing of LiFePO₄/C, acetylene black and PTFE). Furthermore, in order to get the exact capacity of the intrinsic LiFePO₄ particles at charging rate of 600 C, we first charged the battery at 600 C, the initial state of the cathode is $LiFePO_4$) and then measured the capacity under low discharging rate (e.g., 3 C discharge rate) (Figure S4a in Supporting Information). Similarly, to get the capacity at discharging rate of 600 C, we first discharged the battery at 600 C (the initial state of the cathode is FePO₄) and then measured the capacity under low charge rate (e.g., 3 C charge rate) (Figure S4b in Supporting Information). We can see that there are clear voltage platforms in the 3 C discharging and 3 C charging curves, and the capacities at 3 C discharge rate and 3 C

charging rate should reflect the exact capacity of delithiation from the intrinsic LiFePO₄ particles at charging rate of 600 C and lithiation from the FePO₄ particles at discharging rate of 600 C, respectively.

We also found that this ultrahigh rate capability in aqueous electrolyte comes with excellent cycling stability (Figure S6 in Supporting Information), that is, the loss of capacity is only 5% after 400 cycles at 10 C, or <15% after 1000 cycles for 45 nm LiFePO₄ crystal. Additional nanostructure and hierarchy were also introduced to LiFePO4 when a hybrid nanohollow composite with the mean pore sizes of 200 nm was assembled with multiwall carbon nanotubes (MWCNTs) (Figure S7 in Supporting Information), which affords similar CDR and cycling stability in aqueous electrolyte; that is, capacities retained at the rates of 120, 300, and 600 C are 97.6, 84.0, and 62.6 mAh g^{-1} , or 60%, 58%, 50%, and 38% of the theoretical capacity, respectively. Compared with the previously reported CDR in aqueous electrolyte,¹⁶ the higher rates achieved here are most likely due to better optimized carbon coating and higher electronic conductivity. Calculation shows that, at the same electronic conductivity and electrode surface-to-volume ratio, CDR in aqueous media could be higher than in nonaqueous electrolytes nearly by 5 times; the underlying mechanism governing this tremendous difference, however, has not been accounted for.

One intuitive explanation for the higher CDR in aqueous electrolyte would resort to the higher Li⁺-diffusion rate (or conductivity) therein.²⁴ Upon closer examination, however, one would readily rule out this factor because the ionic conductivity in 0.5 M Li₂SO₄ aqueous electrolyte (1.67 mS cm⁻¹) is only higher than that of nonaqueous electrolyte (1.44 mS cm⁻¹ for 1 M LiClO₄ in EC/DMC) by a fraction (Table S4 in Supporting Information). Moreover, the high CDR in aqueous electrolyte (Table S3 in Supporting Information). These evidence suggest that the Li⁺-diffusion in bulk electrolyte should not be the rate-determining step that exercises kinetic control over CDR. This inference is further supported by the numerical analysis on cyclic voltammogram (CV) as shown in Figure 2a and in Figure



Figure 3. Ab initio calculated H_2O/EC absorption at LiFePO₄ and FePO₄ surfaces. (a) Top view of LiFePO₄/vacuum, LiFePO₄/H₂O, and FePO₄/H₂O at (010) direction. The four LiO₆ octahedrons correspond to four Li diffusion channels. The dashed red cycles denote the O vacancies at the surface. (b) to (d) Side view of LiFePO₄/vacuum, LiFePO₄/H₂O, and FePO₄/H₂O at (010) direction, respectively. "HSLI" denotes the hydrated solid–liquid interface. (e) and (f) Side view of LiFePO₄/EC and FePO₄/EC at (010) direction, respectively.

S8a in Supporting Information, where electrochemical kinetics process model (see section S3 in Supporting Information for details) was adopted while taking into account the anodic and cathodic peak shifts upon the change of scanning rate. As visually hinted by the separations between cathodic and anodic peaks, which is much smaller in aqueous (~0.1 V) with minimum overpotential than that in nonaqueous (~ 0.5 V) electrolyte, the overall (surface plus bulk) Li⁺-diffusion constant within LiFePO4 nanoparticles when evaluated in aqueous electrolyte is higher by an order of magnitude than the corresponding values if evaluated in nonaqueous electrolytes, that is, 5.0×10^{-11} cm²/s versus 5.0×10^{-12} cm²/s. Apparently, this gap should be the defining factor for the much higher CDR realized in aqueous-based LIB. Additional evidence comes from the electrochemical impedance spectra, where the interfacial resistance observed for the same nanosized LiFePO44 is much smaller in aqueous electrolyte (Figures S8b and S9b in Supporting Information). Because Li⁺-diffusions inside the nanoparticles of the same source should be independent from the electrolytes used, either aqueous or nonaqueous, the difference has to arise from the corresponding solid/liquid interfaces. This would explain why the modeled overall diffusion constant 5.0 \times $10^{-11}~\text{cm}^2/\text{s}$ is much lower than the theoretical bulk Li+-diffusion constant (1.0 \times 10^{-8} \sim 1.0 \times 10^{-10} cm²/s) within LiFePO₄.²⁶

For the aqueous scenario, the Coulombic efficiency is 100% even at the first cycle, indicating a simple redox process without Li loss. This perhaps should be attributed to the moderate redox potential (3.5 V) of LiFePO₄, which is not high enough to require the formation of an interphase,² differing from mainstream metal oxide cathodes such as $LiCoO_2$, $LiNiO_2$, $LiNiO_2$, $LiNi_xCo_yMn_zO_2$, and $LiMn_2O_4$ that operate beyond 4.0 V.^{27–29} This simple LiFePO₄/electrolyte interface in absence of interphase makes it ideal for ab initio theoretical investigations under the density functional theory (DFT). It should be emphasized here that the carbon coating does not block the electrolyte from directly accessing the LiFePO₄ surfaces.^{29,30} In our model, the (010) surface of LiFePO₄ (Figure 3a and b) is selected as the contact surface with electrolyte, which is normal to the bulk Li diffusion channel.³¹ In such a 2-D classical interface, the inner-Helmholtz structure of LiFePO₄ surfaces constitutes the key conduction passage for Li⁺. The H₂O adsorption at different atomic sites on LiFePO₄ (010) surface, hence, was calculated with varying numbers of H₂O molecules (section S4 in Supporting Information), and the most stable structure is identified when three H₂O molecules adsorb at three different sites as illustrated in Figure 3a and c (the 1, 2, 3 sites). These sites are the exact locations of the O vacancy at the corners of FeO₆ and LiO₆ octahedrons in a stoichiometric (010) LiFePO₄ surface. The O atoms from the adsorbed H_2O molecules fill in such vacancies and complete the originally truncated octahedral coordination symmetry. In particular, at the concave site 1 (Figure 3a and c), the O from H_2O completes both a FeO_6 and a LiO_6 octahedron, with a H_2O -tosurface binding energy being 1.37 eV. The H atom from the same H₂O molecule also forms a strong hydrogen bond with



Figure 4. Experimental evidence for the optimized hydrated interface and the reaction profiles for Li-ions transport across the FePO₄/water interface and FePO₄/EC interface. (a) Total change of mass (*M*) induced by every molar electrons during the CV scan for charge (red line and dots) and discharge (blue dots) process versus the average LiFePO₄ particle size (l = 2r) in 0.5 M Li₂SO₄ aqueous electrolyte. (b) Fourier transform infrared (FTIR) spectra for water, water with LFP, and water with FPO. (c) and (d) Reaction profiles for Li-ions transport across the FePO₄/water interface (c) and FePO₄/EC interface (d) in the discharge process and their energies at each step (right-hand panels). Li, green; Fe, gray; P, purple; O, red; C, brown; H, white.

the neighboring O atom at site 4, pulling this $O_{(site 4)}$ atom toward it by about 0.2 Å (Figure 3a). Note that, beside these three H₂O molecules, no other water molecule has direct access to the LiFePO₄ interface. Such unique arrangement of water molecules at those LiFePO₄ (010) sites actually creates a "Janus" interface in the key Li+-passage that simulates the structures of both solid (electrode) and liquid (electrolyte) phases (Figure 3c). It has been well-established in surface sciences that interfaces of such properties are effective in promoting fast mass transfer.^{32,33} Water adsorption was also calculated at the fully delithiated FePO₄ surfaces. In this case, except the site 1 for the FeO₆ octahedron, the O vacancy sites of LiO₆ octahedron no longer exist. As a result, only one H₂O molecule can be accommodated at the surface at site 1 (Figure 3a and d) with a binding energy of 0.97 eV. On the basis of the above first-principles models, a simple schematic illustration of solid/liquid interfaces of LiFePO₄/water, FePO₄/water, LiFe PO_4/EC , and $FePO_4/EC$ is shown in Figure S16 in Supporting Information.

The transition from "tri-H₂O molecule bound to LiFePO₄ surface" to "mono-H₂O molecule bound to FePO₄ surface" is further verified with an accurate in situ mass measurement, that is, electrochemical quartz crystal microbalance (EQCM). During the charging/discharging process, EQCM in the aqueous electrolyte shows that every molar electron (as determined from the electric current passed) is accompanied by a decrease or increase of 12.6/12.9 g mol⁻¹ mass change (Figure S19 in Supporting Information and red lines in Figure 2b), which is 5.66/5.96 g mol⁻¹ in excess of the deintercalation/intercalation of one Li⁺ (6.94 g mol⁻¹). This change in mass is stable and reversible during the electrochemical charging and discharging. In contrast, mass change per molar electron as registered by EQCM becomes 6.6-6.7 g mol⁻¹ in nonaqueous electrolyte, consistent with the mass

change of one Li⁺ (blue lines in Figure 2b). In a multinanoparticle electrode, it has been well established that one nanoparticle will be fully charged (or discharged) before other nanoparticles start to participate the reaction. Thus, this mass change during the charging/discharging process should correspond to mass change between pure LiFePO₄ and pure FePO₄ nanoparticles, that is, the change from tri-H₂O molecule adsorptions to mono-H₂O adsorption. Using such a model, the excess mass can be calculated while taking into account the particle size and the surface areas (section \$5 in Supporting Information). The calculated excess mass agrees well with the excess mass as measured by EQCM for different nanoparticle sizes as shown in Figure 4a. This is by far the most compelling evidence for the reversible adsorption/desorption of water molecules during the charging/discharging process because there is no other plausible explanation for the measured excess mass change. Fourier transform infrared spectroscopy (FTIR) was also used to study the microstructure of surface H₂O adsorption (section S6 in Supporting Information). The phonon modes related to H-O bond bending and stretching is of particular interest, because this mode at about 3500 cm⁻ does not overlap with other phonon modes. For the LiFePO₄/ H₂O system (red line in Figure 4b), an additional small peak around 2800 cm⁻¹ (compared to the simple summation of LiFePO₄ and H₂O FTIRs) was observed, which was down shifted from the original H-O stretching mode main peak at 3500 cm⁻¹. DFT calculations of the phonon modes using the atomic structure of LiFePO4·3H2O (Table S6 in Supporting Information) also show a soft O(site 1)-H stretching mode of 2500 cm⁻¹ (shifted from the main peak around 3400 cm⁻¹), which arises due to the strong hydrogen bond along the $O_{(site 1)}$ -H- $O_{(site 4)}$ in Figure 3a, while weakening the O-H bond strength. Due to possible complication of other H₂O molecules and the deficiency of using DFT to accurately describe the hydrogen bond, a more quantitative comparison is not possible at the moment. However, the FTIR does qualitatively support our picture of water-binding. Note that in the case of FePO₄·H₂O, both the FTIR and DFT calculation do not show any significant sharp features (Figure 4b and Table S5 in Supporting Information).

The above water absorption picture provides a possible scenario for Li ion diffusion at the solid/liquid interface (section S7 in Supporting Information). At zero or partial intercalation (Li_xFePO₄ where $0 \le x \le 1.0$), the Li⁺ will be accommodated inside the nanocrystal to form a solid solution as proposed by Malik et al.¹² As a result, the surface can be considered as a pure FePO₄ phase and only one water molecule is required to compensate the truncated octahedral symmetry (Figure 3a). Only once full intercalation along any Li channel to generate Li_xFePO₄ phase in this nanoparticle, three H₂O would be required at the surface of the Li channel. On the other hand, it is already known that each Li⁺ in aqueous electrolyte is always coordinated by four water molecules in its primary ^{34–36} In solvation sheath, forming a complex cation $Li^+(H_2O)_4$. order for a Li⁺ to intercalate into the nanoparticle, two but only two water molecules need to be stripped away from this complex (i to ii, and to iii in Figure 4c). Following this initially partial desolvation, the resultant $Li^{+}(H_2O)_2$ can approach the LiFePO₄ surface with almost no additional barrier and docks at the site (iii to iv in Figure 4c) to form a structure similar to the scenario of 3H₂O on top of LiFePO₄ of Figure 3a. After this Li⁺ diffused into the LiFePO₄ bulk along the Li channel, these two H₂O will desorb from the surface. The DFT calculated energies

of these intermediate stages are shown in Figure 4c, and all the major energy barriers were proven to originate from Li⁺ desolvation process at the interface instead of within either electrolyte or electrode bulk. Note that, due to the small size of H₂O molecule, the Li⁺(H₂O)₄ can intimately approach those concaves with truncated symmetries before desolvation occurs, making it possible for the surface electric field to facilitate the desolvation process, and hence further reduce the energy barriers of the discharging process from the ones shown in Figure 4c.

Following the same procedure, the possible interaction between nonaqueous solvent EC molecule and LiFePO₄ (010) surface is also evaluated. Previously, it has been found that EC is the preferred solvent member in the primary solvation sheath of Li⁺ despite the dominant DMC presence.³⁷ It was found that for both LiFePO₄ and FePO₄, only one EC molecule can be adsorbed due to steric hindrance (Figure 3e and f), and three EC molecules need to be desolvated from the complex cation $Li^{+}(EC)_{4}^{38,39}$ before it can approach the concaves at (010) surface in order for the intercalation of one Li⁺ to happen. This not-so-intimate interaction fails to compensate the truncated solid symmetry at those (010) channels, and Li⁺ must overcome a higher energy barriers to cross such interfaces, as shown in Figure 4d and consistent with earlier electrochemical experiments, where the height of the barrier was estimated to be as high as 50 kJ/mol.⁴⁰ Likewise, the desolvation process in nonaqueous electrolytes has to happen far away from the surface due to the large EC molecule size, which makes its desolvation less likely to be assisted by the surface electric field. Combination of all these factors led to the much higher barrier for Li⁺ to cross such a solid/liquid interface, resulting in much lower CDR. We calculated the potential barriers for the charging process using the same approach, and the results are shown in Figure S24 in Supporting Information. Similarly, we saw a lower barrier for the aqueous electrolyte than the nonaqueous electrolyte thanks to the optimized interface created by the water adsorption.

Conclusions. In summary, Janus solid-liquid interface was created by the unique arrangement of water molecules at the specific (010) sites of LiFePO₄, whose partial-solid and partialliquid structure compensates the truncated surface symmetry of the former, and serves as a transition state during the Li⁺transport across the electrolyte/electrode interface. This coordination of Li⁺ by O_s and O_L from solid electrode and liquid electrolyte respectively relieves the energy barriers associated with the Li⁺-desolvation process, leading to minimum overpotential and a least-resisting interface that are responsible for the ultrahigh CDR observed. Energy storage devices constructed upon such nanosized LiFePO₄ in aqueous electrolytes are expected to demonstrate both unprecedented rate capability and maximum round trip efficiency for energy conversion during the operation. Of more fundamental and broader significance, the above picture of $LiFePO_4$ (010) surface in water actually represents what an ideal interface shouldbe it adopts a transition structure (symmetry) that closely imitates the interior bulk of both phases so that a mobile species would experience the minimum discontinuities in properties while transporting across such interfaces. We believe that this should be the principle that guides the rational design of a perfect interface of ultrafast transport properties.

Methods Section. *Experimental Section.* The LiFePO₄ particles with mean size of 45 nm (calculated from Scherrer equation) were synthesized by reflux route in ethylene glycol

solution under atmospheric pressure. In a typical route, FeSO₄· $7H_2O_1$, H_3PO_4 , and LiOH·H₂O were used as starting materials in a molar ratio of 1:1.5:2.7 and ethylene glycol (EG) was applied as solvent. H₃PO₄ was slowly introduced to the LiOH solution under stirring, and then FeSO₄ solution was added into the mixture. After homogenization, the reaction mixture was heated at reflux condition for 10 h under Ar atmosphere. The resulting suspension was washed several times with water and ethanol. To achieve carbon coating, LiFePO₄ nanoparticles were mixed with 20 wt % of glucose and then carbonized at 650 °C for 6 h in Ar atmosphere. For the synthesis of LiFePO₄ nanoparticles with mean size of 101 nm (calculated from Scherrer equation), the feeding sequence was changed. In this case, H₃PO₄ was slowly introduced to the FeSO₄ solution under stirring. Then LiOH solution was added into the mixture. Other steps are the same as the synthetic process for 45 nm LiFePO₄ particles.

The working electrode was prepared by pressing a mushy mixture of the prepared composite, acetylene black and poly(tetrafluoroethylene) (PTFE) (weight ratio: 30:50:20) dispersed in isopropanol onto steel grid. The thickness of the electrode was about 50 μ m, and its area was about 0.3 cm². The active materials mass loading is 0.8 mg. The electrode area was about 0.3 cm² and its weight was about 0.8 mg. Aqueous Li_2SO_4 electrolyte (0.5 M) was used as the electrolyte. Galvanostatic charge-discharge tests (Maccor, MC-16 Battery Test System) of the LiFePO $_4$ /C composites electrodes were performed using a three-electrode cell, where the active carbon and saturated calomel electrode (SCE) were used as counter and reference electrodes, respectively. Pt wire and Ag/AgCl were used as the counter and reference electrodes in cyclic voltammetry (CV) and electrochemical impedance (CHI 660E, Shanghai Chenhua). The electrochemical impedance spectra were recorded from 10^4 to 0.1 Hz and the amplitude of the used perturbation was 10 mV. For the EQCM test on an electrochemical workstation (CHI 440B, Shanghai Chenhua), the LiFePO₄ particles were deposited on the gold disk which was coated on the crystal plate. Ionic conductivities of the electrolyte solutions were measured at room temperature using a Model DDS-307 conductometer (Shanghai REX Instrument Factory, China).

Theoretical. All calculations are performed using the planewave projector-augmented wave method⁴¹ with an energy cutoff of 520 eV, as implemented in the Vienna ab initio simulation package (VASP).⁴²⁻⁴⁶ The Perdew-Burke-Ernzerhof (PBE)⁴⁷ form of generalized gradient approximation (GGA) is chosen as the exchange-correlation potential. The PBE+U approach⁴⁸ is employed to take account of the strong on-site Coulomb interaction (U) presented in the localized 3d electrons of Fe, with the U values set to 4.7 eV in LiFePO₄ and 5.9 eV in FePO₄. To obtain reliable optimized structures and the total energy, all the atomic positions and cell parameters are fully relaxed using a conjugate gradient algorithm, until the force on each atom is smaller than 0.02 eV/Å and energies are converged to within 5×10^{-6} eV per atom. A $8 \times 4 \times 1$ k-point grid within the Monkhorst-Pack scheme⁴⁹ is used to sample the Brillouin zone of the LiFePO₄ (010) surface unit cell, and a $1 \times 1 \times 1$ k-point grid is used for calculations of separate molecules. We finally employed a climbing-image nudged elastic band (cNEB) method⁵⁰ to calculate the energy barriers for Li atoms diffusion in the bulk LiFePO₄ (FePO₄) and from the subsurface to the surface. A DFT-D₂ semiempirical dispersion-correction approach⁵¹ to correct the van der Waals

(vdW) interactions and a dipole correction method are also tested in our calculations, and we find there is little influence to the final optimized structures. A ferromagnetic high-spin Fe state is assumed, and the energetic effects of the magnetic ordering are small. In order to study the interfacial properties between LiFePO₄/FePO₄ surfaces and electrolytes, we build a 1 $\times 2 (a \times c)$ bilayer LiFePO₄/FePO₄ (010) slab to simulate the $LiFePO_4/FePO_4$ (010) surface. A vacuum buffer space of at least 12 Å is set for slabs with adsorbed molecules. Lattice parameters of the supercell (including slab and vacuum) are fixed, and the inner part of the slab is frozen at the bulk positions and only the atoms near the surface are allowed to relax. The electron counting rule^{52,53} and the surface reconstruction are also considered. The vibrational modes are calculated based on the numerical calculations of the second derivatives of the potential energy within the harmonic approximation.

Additional details on materials and methods are included in the Supporting Information.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.nano-lett.5b02379.

Additional materials and methods; additional materials characterization and electrochemical performances; electrochemical kinetics analysis; ab initio calculation results; the mass change for LiFePO₄ during charging and discharging process versus the average LiFePO₄ particle size; Fourier transform infrared (FTIR) spectra observation and vibrational modes calculations for the interface structure; detailed discussions for the discharge and charge process. (PDF)

AUTHOR INFORMATION

Corresponding Authors

*E-mail: wuyp@fudan.edu.cn.

*E-mail: panfeng@pkusz.edu.cn.

Author Contributions

These authors contributed equally to this work (J.Z., Y.H., Y.D).

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENTS

The research was financially supported by National Project for EV Batteries (20121110, OptimumNano, Shenzhen), National Distinguished Young Scientists of China (51425301), STCSM (12JC1401200), Guangdong Innovation Team Project (No. 2013N080), and Shenzhen Science and Technology Research Grant (No. ZDSY20130331145131323 and CXZZ20120829172325895). L.W.W. is supported through the Theory of Material project by the Director, Office of Science (SC), Basic Energy Science (BES)/Materials Science and Engineering Division (MSED) of the U.S. Department of Energy (DOE) under the contract No. DE-AC02-05CH11231. Additionally, we acknowledge the support of ShenZhen National Super Computing Center.

REFERENCES

(1) Xu, K. Chem. Rev. 2004, 104, 4303-4418.

Nano Letters

- (2) Xu, K. Chem. Rev. 2014, 114, 11503-11618.
- (3) Jow, T. R.; Xu, K.; Borodin, O.; Ue, M. Electrolytes for lithium and lithium-ion batteries; Springer: New York, 2014; Vol. 58.
- (4) Dunn, B.; Kamath, H.; Tarascon, J.-M. Science 2011, 334, 928–935.
- (5) Yang, Z. G.; Zhang, J.; Kintner-Meyer, M. C.; Lu, X. C.; Choi, D.; Lemmon, J. P.; Liu, J. *Chem. Rev.* **2011**, *111*, 3577–3613.
- (6) Presser, V.; Dennison, C. R.; Campos, J.; Knehr, K. W.; Kumbur, E. C.; Gogotsi, Y. Adv. Energy Mater. **2012**, *2*, 895–902.
- (7) Halls, J. E.; Hawthornthwaite, A.; Hepworth, R. J.; Roberts, N. A.;
- Wright, K. J.; Zhou, Y.; Haswell, S. J.; Haywood, S. K.; Kelly, S. M.; Lawrence, N. S.; Wadhawan, J. D. Energy Environ. Sci. 2013, 6, 1026-
- 1041.
- (8) Arico, A. S.; Bruce, P.; Scrosati, B.; Tarascon, J.-M.; van Schalkwijk, W. Nat. Mater. 2005, 4, 366–377.
- (9) Whittingham, M. S. Chem. Rev. 2014, 114, 11414-11443.
- (10) Gaberscek, M.; Kuzma, M.; Jamnik, J. Phys. Chem. Chem. Phys. 2007, 9, 1815–1820.
- (11) Liu, J.; Kunz, M.; Chen, K.; Tamura, N.; Richardson, T. J. J. Phys. Chem. Lett. 2010, 1, 2120–2123.
- (12) Malik, R.; Abdellahi, A.; Ceder, G. J. Electrochem. Soc. 2013, 160, A3179–A3197.
- (13) Ravet, N.; Chouinard, Y.; Magnan, J. F.; Besner, S.; Gauthier, M.; Armand, M. J. Power Sources **2001**, 97–98, 503–507.
- (14) Gaberscek, M.; Dominko, R.; Jamnik, J. Electrochem. Commun. 2007, 9, 2778–2783.
- (15) Jansen, A. N.; Dees, D. W.; Abraham, D. P.; Amine, K.; Henriksen, G. L. J. Power Sources **2007**, 174, 373–379.
- (16) He, P.; Zhang, X.; Wang, Y. G.; Cheng, L.; Xia, Y. Y. J. Electrochem. Soc. 2008, 155, A144-A150.
- (17) Dathar, G. K. P.; Sheppard, D.; Stevenson, K. J.; Henkelman, G. *Chem. Mater.* **2011**, *23*, 4032–4037.
- (18) Wang, Q.; Zakeeruddin, S. M.; Exnar, I.; Grätzel, M. J. Electrochem. Soc. 2004, 151, A1598–A1603.
- (19) Abe, T.; Sagane, F.; Ohtsuka, M.; Iriyama, Y.; Ogumi, Z. J. Electrochem. Soc. 2005, 152, A2151-A2154.
- (20) Li, W.; Dahn, J.; Wainwright, D. Science 1994, 264, 1115–1118.
 (21) Luo, J.-Y.; Cui, W.-J.; He, P.; Xia, Y.-Y. Nat. Chem. 2010, 2, 760–765.
- (22) Tang, W.; Liu, L.; Zhu, Y.; Sun, H.; Wu, Y.; Zhu, K. Energy Environ. Sci. 2012, 5, 6909-6913.
- (23) Tang, W.; Hou, Y.; Wang, F.; Liu, L.; Wu, Y.; Zhu, K. Nano Lett. **2013**, *13*, 2036–2040.
- (24) Tang, W.; Zhu, Y. S.; Hou, Y. Y.; Liu, L.; Wu, Y. P.; Loh, K. P.; Zhang, H. p.; Zhu, K. Energy Environ. Sci. 2013, 6, 2093–2104.
- (25) Wang, X. J.; Hou, Y. Y.; Zhu, Y. S.; Wu, Y. P.; Holze, R. Sci. Rep. **2013**, 3, 1401.
- (26) Malik, R.; Burch, D.; Bazant, M.; Ceder, G. Nano Lett. 2010, 10, 4123-4127.
- (27) Edström, K.; Gustafsson, T.; Thomas, J. O. Electrochim. Acta 2004, 50, 397-403.
- (28) Dupré, N.; Martin, J.-F.; Degryse, J.; Fernandez, V.; Soudan, P.; Guyomard, D. J. Power Sources **2010**, *195*, 7415–7425.
- (29) Akita, Y.; Segawa, M.; Munakata, H.; Kanamura, K. J. Power Sources **2013**, 239, 175–180.
- (30) Zaghib, K.; Dontigny, M.; Charest, P.; Labrecque, J. F.; Guerfi, A.; Kopec, M.; Mauger, A.; Gendron, F.; Julien, C. M. *J. Power Sources* **2008**, *185*, 698–710.
- (31) Wang, L.; Zhou, F.; Meng, Y. S.; Ceder, G. Phys. Rev. B 2007, 76, 165435.
- (32) de Gennes, P. G. Angew. Chem., Int. Ed. Engl. 1992, 31, 842– 845.
- (33) Hu, J.; Zhou, S.; Sun, Y.; Fang, X.; Wu, L. Chem. Soc. Rev. 2012, 41, 4356–4378.
- (34) Rudolph, W.; Brooker, M. H.; Pye, C. C. J. Phys. Chem. 1995, 99, 3793-3797.
- (35) Pye, C. C.; Rudolph, W.; Poirier, R. A. J. Phys. Chem. **1996**, 100, 601–605.
- (36) Mähler, J.; Persson, I. Inorg. Chem. 2011, 51, 425-438.

- (37) von Cresce, A.; Xu, K. Electrochem. Solid-State Lett. 2011, 14, A154–A156.
- (38) Yanase, S.; Oi, T. J. Nucl. Sci. Technol. 2002, 39, 1060-1064.
- (39) Bogle, X.; Vazquez, R.; Greenbaum, S.; Cresce, A. v. W.; Xu, K. J. Phys. Chem. Lett. **2013**, *4*, 1664–1668.
- (40) Xu, K.; von Cresce, A.; Lee, U. Langmuir 2010, 26, 11538-11543.
- (41) Blöchl, P. Phys. Rev. B: Condens. Matter Mater. Phys. 1994, 50, 17953-17979.
- (42) Kresse, G.; Furthmüller, J. Comput. Mater. Sci. 1996, 6, 15-50.
- (43) Kresse, G.; Furthmüller, J. Phys. Rev. B: Condens. Matter Mater. Phys. 1996, 54, 11169–11186.
- (44) Kresse, G.; Hafner, J. Phys. Rev. B: Condens. Matter Mater. Phys. 1993, 47, 558–561.
- (45) Kresse, G.; Hafner, J. Phys. Rev. B: Condens. Matter Mater. Phys. 1994, 49, 14251–14269.
- (46) Kresse, G.; Joubert, D. Phys. Rev. B: Condens. Matter Mater. Phys. 1999, 59, 1758-1775.
- (47) Perdew, J.; Burke, K.; Ernzerhof, M. Phys. Rev. Lett. 1996, 77, 3865–3868.
- (48) Anisimov, V.; Zaanen, J.; Andersen, O. *Phys. Rev. B: Condens. Matter Mater. Phys.* **1991**, *44*, 943–954.
- (49) Monkhorst, H.; Pack, J. Phys. Rev. B 1976, 13, 5188-5192.
- (50) Henkelman, G.; Uberuaga, B. P.; Jónsson, H. J. Chem. Phys.
- **2000**, *113*, 9901–9904. (51) Tkatchenko, A.; Scheffler, M. Phys. Rev. Lett. **2009**, *102*, 073005.
- (52) Tasker, P. J. Phys. C: Solid State Phys. **1979**, 12, 4977–4984.
- (53) Manna, L.; Wang; Cingolani, R.; Alivisatos, A. P. J. Phys. Chem. B 2005, 109, 6183-6192.